

# Interlayer Engineering of $\alpha$ -MoO<sub>3</sub> Modulates Selective Hydronium Intercalation in Neutral Aqueous Electrolyte

Haozhe Zhang, Weixing Wu, Qiyu Liu, Fan Yang, Xin Shi, Xiaoqing Liu, Minghao Yu,\* and Xihong Lu\*

**Abstract:** Among various charge-carrier ions for aqueous batteries, non-metal hydronium ( $H_3O^+$ ) with small ionic size and fast diffusion kinetics empowers  $H_3O^+$ -intercalation electrodes with high rate performance and fast-charging capability. However, pure  $H_3O^+$  charge carriers for inorganic electrode materials have only been observed in corrosive acidic electrolytes, rather than in mild neutral electrolytes. Herein, we report how selective  $H_3O^+$  intercalation in a neutral  $ZnCl_2$  electrolyte can be achieved for water-proton co-intercalated  $\alpha$ -MoO<sub>3</sub> (denoted WP-MoO<sub>3</sub>).  $H_2O$  molecules located between MoO<sub>3</sub> interlayers block  $Zn^{2+}$  intercalation pathways while allowing smooth  $H_3O^+$  intercalation/diffusion through a Grotthuss proton-conduction mechanism. Compared to  $\alpha$ -MoO<sub>3</sub> with a  $Zn^{2+}$ -intercalation mechanism, WP-MoO<sub>3</sub> delivers the substantially enhanced specific capacity (356.8 vs. 184.0  $mA\text{h g}^{-1}$ ), rate capability (77.5 % vs. 42.2 % from 0.4 to 4.8  $Ag^{-1}$ ), and cycling stability (83 % vs. 13 % over 1000 cycles). This work demonstrates the possibility of modulating electrochemical intercalating ions by interlayer engineering, to construct high-rate and long-life electrodes for aqueous batteries.

## Introduction

Rechargeable aqueous batteries with neutral electrolytes have attracted intensive scientific attention as promising alternatives for large-scale energy storage technologies. The utilized water-based electrolytes offer significant advantages

of high ionic conductivity ( $\approx 1 \text{ S cm}^{-1}$ ), simplified manufacture, low cost, and intrinsic safety.<sup>[1]</sup> In particular, Zn metal batteries (ZMBs) with mild aqueous electrolytes have recently stood out, due to the direct use of Zn metal anodes with a high specific capacity ( $\approx 820 \text{ mA h g}^{-1}$ ) and a low stripping/plating potential ( $-0.76 \text{ V}$  vs. standard hydrogen electrode).<sup>[2]</sup> Numerous efforts have been devoted to exploring  $Zn^{2+}$ -host cathode materials of ZMBs, which has brought Mn-compounds, V-compounds, Prussian blue materials into the spotlight.<sup>[3]</sup> However, the  $Zn^{2+}$ -intercalation chemistry generally shows sluggish kinetics and unsatisfactory cycling stability.<sup>[4]</sup> In aqueous electrolytes,  $Zn^{2+}$  tends to form a large-size hydrated state ( $Zn(H_2O)_6^{2+}$  with 5.5 Å) due to the strong  $Zn^{2+}$ -water interaction.<sup>[5]</sup> The intercalation of  $Zn^{2+}$  into cathode hosts thus requires large de-solvation and intercalation energy. Besides, bivalent  $Zn^{2+}$  imposes a strong repulsive force with the hosts, leading to the large  $Zn^{2+}$ -diffusion energy barriers within the hosts and the undesired structure distortion of hosts.<sup>[6]</sup>

Apart from  $Zn^{2+}$ , non-metal hydronium ( $H_3O^+$ ) has also been recognized as favorable charge carrier ions for aqueous batteries. Assigned to the small ionic size ( $\approx 1.0 \text{ \AA}$ ) and light molecular mass,  $H_3O^+$  intercalation presents attractive high-kinetics and highly reversible behaviors.<sup>[7]</sup> The partial involvement of  $H_3O^+$  intercalation was also discovered for the charge-storage mechanism of ZMB cathodes.<sup>[8]</sup> For example, Sun et al. uncovered the consequent intercalation of  $H_3O^+$  and  $Zn^{2+}$  for  $\epsilon$ -MnO<sub>2</sub> cathode in a mixed  $ZnSO_4/MnSO_4$  electrolyte.<sup>[9]</sup> The charge-transfer resistance of  $\epsilon$ -MnO<sub>2</sub> in the  $H_3O^+$ -intercalation step is three orders of magnitude smaller than that in the  $Zn^{2+}$ -intercalation step. A similar phenomenon was also observed for polyaniline-intercalated MnO<sub>2</sub> nanolayers, in which the diffusion coefficient of  $H_3O^+$  ( $5.84 \times 10^{-12} \text{ cm}^2 \text{ s}^{-1}$ ) was substantially higher than that of  $Zn^{2+}$  ( $7.35 \times 10^{-14} \text{ cm}^2 \text{ s}^{-1}$ ).<sup>[10]</sup> These findings inspire that selective  $H_3O^+$  intercalation into cathodes would bring the constructed ZMBs with significant performance advance in terms of capacity, kinetics, as well as cycle life. However, thus far, pure  $H_3O^+$ -intercalation behavior for layered/tunneled cathodes has only been observed in corrosive acidic electrolytes.<sup>[7b,11]</sup> It remains a grand challenge to achieve selective  $H_3O^+$  intercalation in mild neutral electrolytes.

In this study, we, taking orthorhombic MoO<sub>3</sub> ( $\alpha$ -MoO<sub>3</sub>) as an example, for the first time demonstrate the feasibility of selective  $H_3O^+$  intercalation in a neutral  $ZnCl_2$  electrolyte.  $\alpha$ -MoO<sub>3</sub> is selected due to its typical layered structure with distorted [MoO<sub>6</sub>] octahedra bilayers weakly bonded by van der Waals force.<sup>[12]</sup> The complete redox of Mo<sup>4+</sup>/Mo<sup>6+</sup> allows

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Figure 1. Schematic illustration of the  $Zn^{2+}$ -intercalation chemistry for  $\alpha\text{-MoO}_3$  and the selective  $H_3O^+$ -intercalation chemistry for WP- $\text{MoO}_3$ .

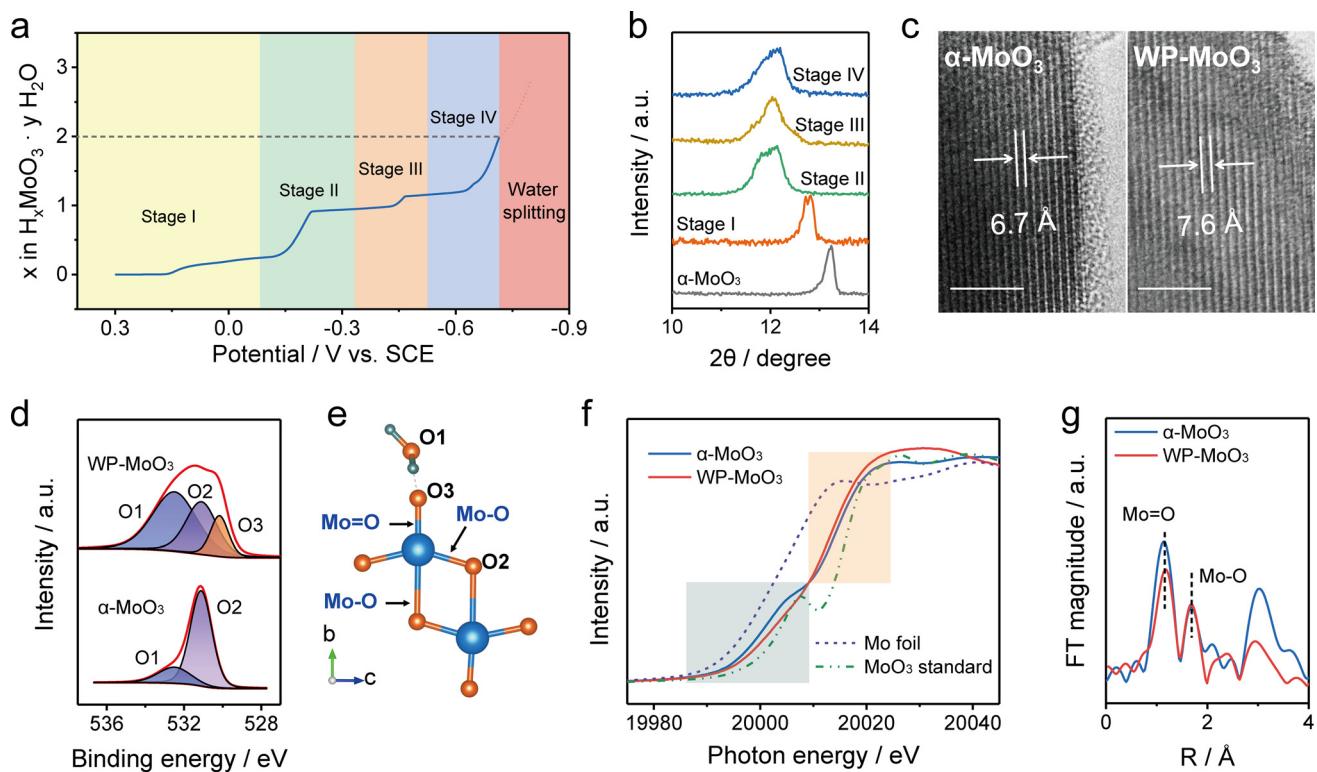
$\alpha\text{-MoO}_3$  with an attractive theoretical capacity of  $372\text{ mA h g}^{-1}$ . Selective  $H_3O^+$ -intercalation chemistry is modulated for  $\alpha\text{-MoO}_3$  through a water-proton co-intercalation strategy (denoted WP- $\text{MoO}_3$ ), which further tackles the low-capacity, poor-rate, and short-life issues faced by pristine  $\alpha\text{-MoO}_3$  with  $Zn^{2+}$ -intercalation chemistry.  $H_2O$  molecules located between WP- $\text{MoO}_3$  interlayers impose a high  $Zn^{2+}$ -intercalation energy barrier by blocking  $Zn^{2+}$  diffusion pathways (Figure 1). Meanwhile,  $H_3O^+$  intercalation/diffusion can be smoothly achieved within WP- $\text{MoO}_3$  interlayers through a well-known Grothuss mechanism (proton jumping between water molecules).<sup>[13]</sup> In contrast to  $Zn^{2+}$ -intercalation  $\alpha\text{-MoO}_3$ , selective  $H_3O^+$ -intercalation WP- $\text{MoO}_3$  depicts substantially enhanced redox depth (1.92 vs. 0.99  $e^-$  per Mo atom; 357 vs. 184  $\text{mA h g}^{-1}$ ), rate capability (77.5% vs. 42.2% from 0.4 to 4.8  $\text{A g}^{-1}$ ), and cycling stability (83% vs. 13% over 1000 cycles).

## Results and Discussion

$\alpha\text{-MoO}_3$  nanoparticles (Figure S1) were first synthesized through a sol-gel method. WP- $\text{MoO}_3$  electrode was obtained from  $\alpha\text{-MoO}_3$  electrode through a controllable and time-efficient ( $\approx 6.1$  min) electrochemical linear sweep voltammetry (LSV) method (Figure S2) in a three-electrode cell with an electrolyte of 1 M  $H_2SO_4$ . Compared with  $\alpha\text{-MoO}_3$  electrode, WP- $\text{MoO}_3$  displays an apparent color change from dark grey to purple (Figure S3), which is attributed to the formation of  $Mo^{4+}/Mo^{5+}$  (Figure S4).<sup>[14]</sup> Almost no morphological variation was observed between  $\alpha\text{-MoO}_3$  and WP- $\text{MoO}_3$ . The amount of intercalated  $H^+$  can be estimated by calculating the amount of charge transfer (Figure 2a). The overall intercalation process presents four stages, referring to potential windows (vs. saturated calomel electrode (SCE)) of 0.3–0.1 V (Stage I), –0.1––0.34 V (Stage II), –0.34––0.53 V (Stage III), and –0.53––0.72 V (Stage IV). Approximately, the intercalated  $H^+$  numbers per  $\text{MoO}_3$  unit are 0.25, 0.75, 0.25, and 0.75 at Stage I, Stage II, Stage III, and Stage IV, respectively. X-ray diffraction (XRD) spectra uncover that the peak position corresponding to the interlayer spacing of  $\alpha\text{-MoO}_3$  gradually shifts towards negative at Stage I and II and keeps almost unchanged at Stage III and IV (Figure 2b &

S5). This peak of WP- $\text{MoO}_3$  is located at  $12.0^\circ$ , indicating the interlayer distance expansion from  $6.7\text{ \AA}$  for  $\alpha\text{-MoO}_3$  to  $7.6\text{ \AA}$  (Figure S6). Besides, new peaks at  $32.5^\circ$ ,  $35.0^\circ$ , and  $37.1^\circ$  are observed for WP- $\text{MoO}_3$ , which indicates a monoclinic phase of proton-intercalated  $\text{MoO}_3$ .<sup>[15]</sup> The widened interlayer distance of WP- $\text{MoO}_3$  was also evidenced by high-resolution transmission electron microscopy (HRTEM) images, in which the interlayer spacings are determined to be  $6.7\text{ \AA}$  and  $7.6\text{ \AA}$  for  $\text{MoO}_3$  and WP- $\text{MoO}_3$ , respectively (Figure 2c & S7).

Figure 2d compares the O 1s X-ray photoelectron spectroscopy (XPS) spectra of  $\text{MoO}_3$  and WP- $\text{MoO}_3$  electrodes. Two peaks located at 532.6 eV and 531.2 eV are observed for  $\alpha\text{-MoO}_3$ , which correspond to lattice O in  $\text{MoO}_3$  (denoted O2) and O in adsorbed  $H_2O$  (denoted O1).<sup>[16]</sup> Notably, WP- $\text{MoO}_3$  shows an O1 peak with the substantially enhanced intensity, verifying the intercalation of  $H_2O$  into WP- $\text{MoO}_3$ . Thermogravimetric analysis (TGA, Figure S8) results of both electrodes also suggest the intercalated  $H_2O$  molecules in WP- $\text{MoO}_3$ . More interestingly, an additional XPS peak at 530.2 eV (denoted O3) is observed for WP- $\text{MoO}_3$ . This peak can be assigned to the terminal O of  $[\text{MoO}_6]$  bilayers, which splits from O2 peak due to the formation of the hydrogen bond with the intercalated  $H_2O/H_3O^+$  (as illustrated in Figure 2e). Furthermore, synchrotron-based X-ray absorption near-edge spectra (XANES) measurements were performed to investigate the localized coordination environments of Mo sites in  $\alpha\text{-MoO}_3$  and WP- $\text{MoO}_3$ . The Mo K-edge XANES spectra of  $\alpha\text{-MoO}_3$  and WP- $\text{MoO}_3$ , as well as standard Mo foil and  $\text{MoO}_3$  as references, are displayed in Figure 2f. In comparison with  $\alpha\text{-MoO}_3$ , the slightly negative-shifted rising-edge of WP- $\text{MoO}_3$  around 20015 eV suggests the enriched electron densities around Mo sites.<sup>[17]</sup> This result is consistent with the analysis of O K-edge XANES spectra, which witness the decreased peak intensity of WP- $\text{MoO}_3$  at the energy region of 530–540 eV (Figure S9).<sup>[18]</sup> In addition, the Mo pre-edge of WP- $\text{MoO}_3$  around 20007 eV, referring to the O 1s–Mo 4d electron transfer, is obviously decreased compared with that of  $\alpha\text{-MoO}_3$ , reflecting the interaction between the terminal O of  $[\text{MoO}_6]$  bilayers and the intercalated species (i.e.,  $H_2O$  and  $H_3O^+$ ).<sup>[19]</sup> To acquire the detailed bonding and coordination information, corresponding R space curves after  $k^2[\chi(k)]$ -weighted Fourier transform of the extended X-ray absorption fine structure (EXAFS) and



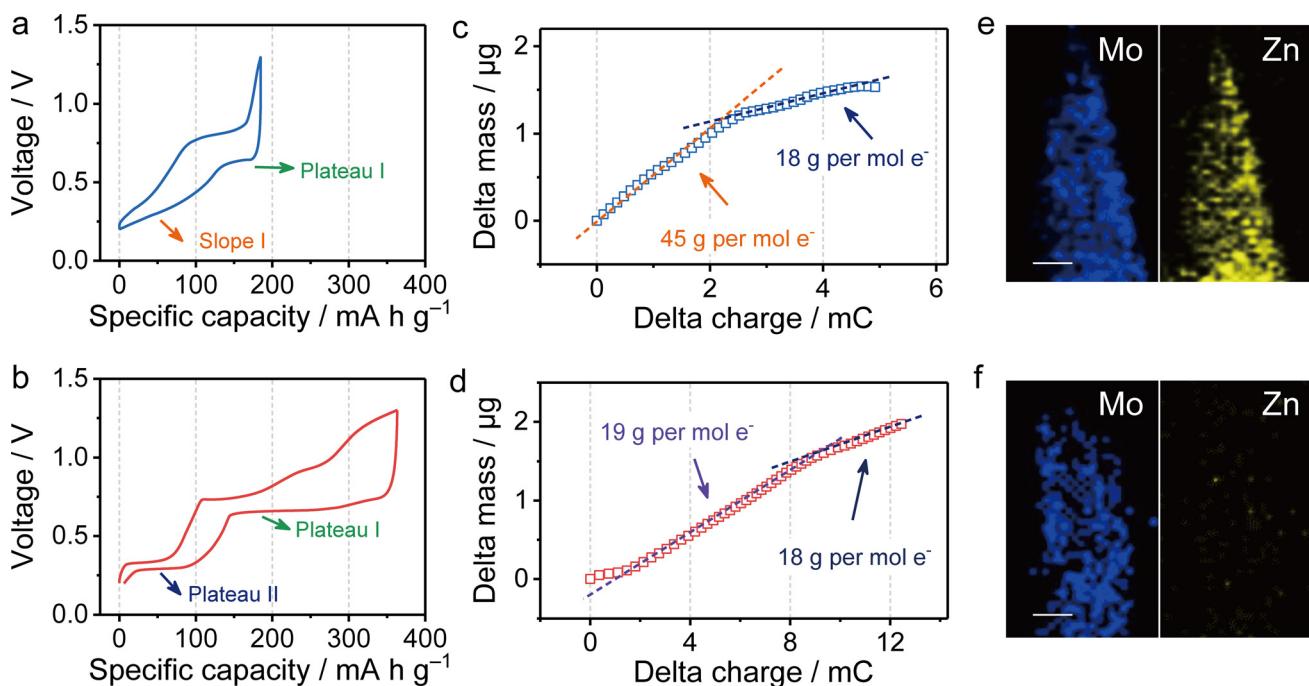
**Figure 2.** a) The intercalated proton amount as a function of potential during the water-proton co-intercalation process. b) XRD patterns of  $\alpha$ - $\text{MoO}_3$  electrode at different intercalation stages. c) HRTEM images of  $\alpha$ - $\text{MoO}_3$  and WP- $\text{MoO}_3$ . Scale bars: 5 nm. d) O 1s XPS spectra of  $\alpha$ - $\text{MoO}_3$  and WP- $\text{MoO}_3$ . e) Illustration of different O atoms in WP- $\text{MoO}_3$ . f) Normalized Mo K-edge XANES spectra of  $\text{MoO}_3$  standard, Mo foil standard,  $\alpha$ - $\text{MoO}_3$ , and WP- $\text{MoO}_3$ . g) Radial distribution functions of  $\alpha$ - $\text{MoO}_3$  and WP- $\text{MoO}_3$  obtained from the  $k^2\chi(k)$  by Fourier transform.

quantitatively fitting spectra are presented in Figure 2g & S10. Two representative peaks at 1.16 and 1.69 Å can be assigned to the scattering of  $\text{Mo}=\text{O}$  and  $\text{Mo}-\text{O}$  bonds in  $[\text{MoO}_6]$  bilayers, respectively.<sup>[20]</sup> These two peaks are also indicated in Figure 2e. Both  $\text{Mo}=\text{O}$  and  $\text{Mo}-\text{O}$  bonds show negligible length difference between  $\alpha$ - $\text{MoO}_3$  (1.72 & 1.96 Å) and WP- $\text{MoO}_3$  (1.73 & 1.97 Å), implying the water-proton pre-intercalation causes minor distortion of octahedron  $[\text{MoO}_6]$  layers. Moreover, the coordination number of  $\text{Mo}-\text{O}$  (from 1.8 to 1.4) and  $\text{Mo}=\text{O}$  (from 2.0 to 1.4) decreases due to the water-proton co-intercalation, which again verifies the interaction between the terminal O of  $[\text{MoO}_6]$  bilayers and the intercalated  $\text{H}_2\text{O}/\text{H}_3\text{O}^+$ .

The ion-intercalation behaviors of  $\alpha$ - $\text{MoO}_3$  and WP- $\text{MoO}_3$  electrodes were explored in two-electrode cells with Zn foil as the anode and 2 M  $\text{ZnCl}_2$  aqueous solution as the electrolyte. Figure 3a and Figure 3b present the galvanostatic charge/discharge (GCD) curves of  $\alpha$ - $\text{MoO}_3$  and WP- $\text{MoO}_3$  electrodes after three-cycle activation at  $0.4 \text{ A g}^{-1}$ . Both electrodes show two mainly ion-intercalation stages in their discharge curves, which agrees well with the cyclic voltammetry (CV) curves (Figure S11). In detail,  $\alpha$ - $\text{MoO}_3$  electrode presents a discharge curve consisting of one plateau region (Plateau I,  $\approx 0.62 \text{ V}$ ) and one slope region (Slope I, 0.25–0.35 V), while WP- $\text{MoO}_3$  electrode presents two plateaus around 0.68 V (Plateau I) and 0.30 V (Plateau II). Impressively, WP- $\text{MoO}_3$  electrode exhibits a high redox depth of  $1.92 \text{ e}^-$  per Mo atom, close to the full conversion of  $\text{Mo}^{4+}/$

$\text{Mo}^{6+}$ . By contrast, a shallow redox depth of  $0.99 \text{ e}^-$  per Mo atom is achieved by  $\alpha$ - $\text{MoO}_3$  electrode.

Electrochemical quartz crystal microbalance (EQCM) was employed to in-operando monitor the mass evolution of  $\text{MoO}_3$  and WP- $\text{MoO}_3$  along with the continuous ion intercalation. In consistent with the discharge curves, the mass change of both electrodes during ion intercalation depicted two stages. At the first stage of  $\alpha$ - $\text{MoO}_3$ , the average weight increase calculated by the curve slope (Figure 3c) is  $45 \text{ g}$  per mol charge, which suggests the co-intercalation of  $\text{Zn}^{2+}$  and  $\text{H}_2\text{O}$  into  $\alpha$ - $\text{MoO}_3$  ( $\text{Zn}^{2+} \cdot 1.5 \text{ H}_2\text{O}$  in average). Meanwhile, the average weight increase of  $\alpha$ - $\text{MoO}_3$  at the second stage is  $18 \text{ g}$  per mol charge, close to the weight of  $\text{H}_3\text{O}^+$  (19 g per mol charge). Clearly,  $\text{Zn}^{2+}/\text{H}_2\text{O}$  co-intercalation plays the dominant role in the charge-storage mechanism of  $\alpha$ - $\text{MoO}_3$ . By contrast, WP- $\text{MoO}_3$  depicts the weight increase of 19 and 18 g per mol charge at the two stages (Figure 3d), verifying the selective  $\text{H}_3\text{O}^+$ -intercalation chemistry of WP- $\text{MoO}_3$  in 2 M  $\text{ZnCl}_2$  electrolyte. This result means although  $\text{Zn}^{2+}$  is the mainly charge carrier in the electrolyte, the charge carrier only consists of  $\text{H}_3\text{O}^+$  in WP- $\text{MoO}_3$  electrode, which comes from the pre-stored  $\text{H}^+$  in WP- $\text{MoO}_3$  and the slight hydrolysis of  $\text{ZnCl}_2$  in the electrolyte. However, the local  $\text{Zn}^{2+}$  concentration near the surface of WP- $\text{MoO}_3$  should be increased during discharging, because  $\text{Zn}^{2+}$  tends to migrate to the Helmholtz layer of WP- $\text{MoO}_3$  due to the electrostatic interaction.



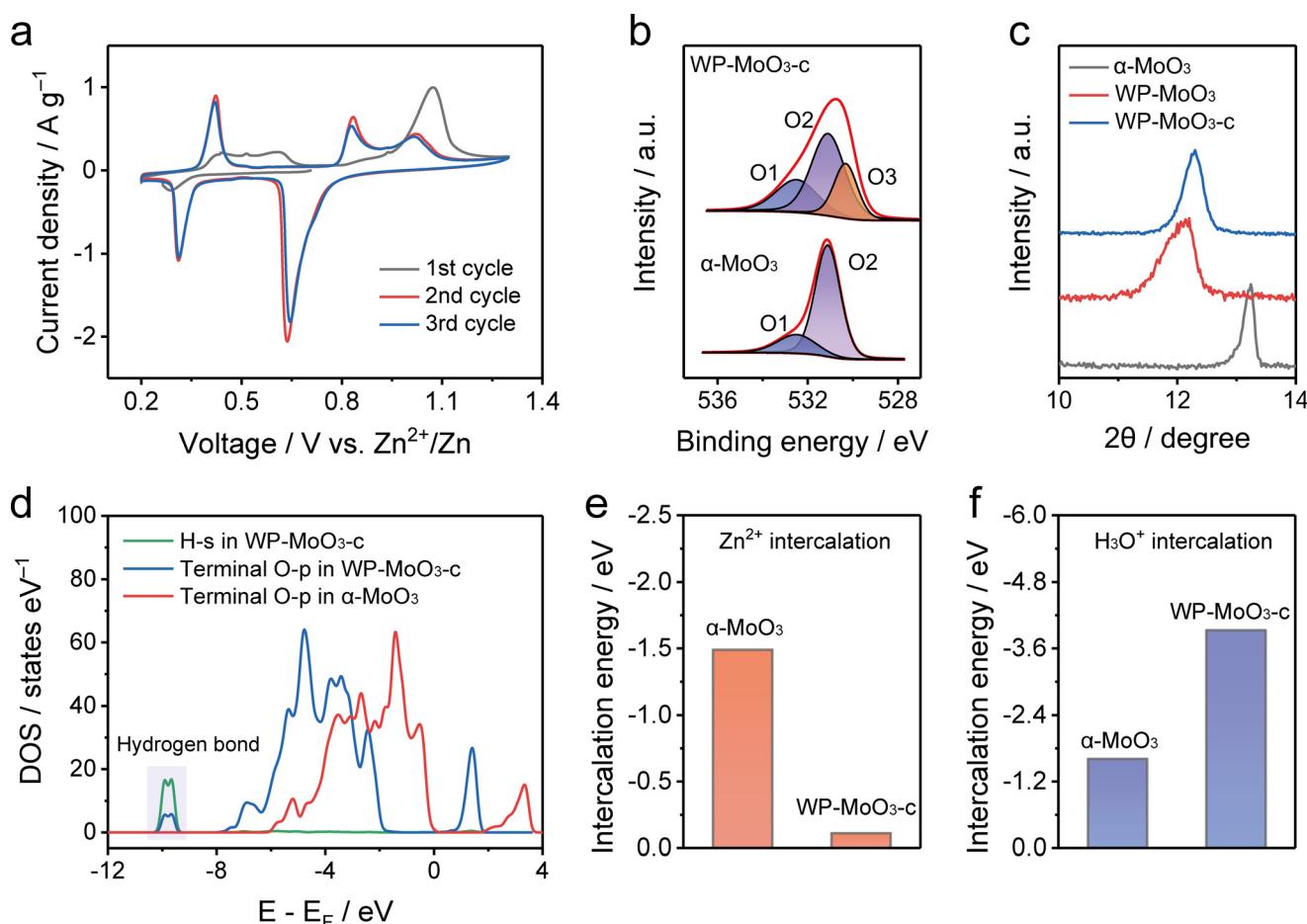
**Figure 3.** GCD profiles of a)  $\alpha$ -MoO<sub>3</sub> and b) WP-MoO<sub>3</sub> electrodes at a current density of  $0.4 \text{ A g}^{-1}$ . Electrode mass change versus charge during the discharge (ion-intercalation) processes of c)  $\alpha$ -MoO<sub>3</sub> and d) WP-MoO<sub>3</sub> electrodes. EDX elemental mapping of e)  $\alpha$ -MoO<sub>3</sub> and f) WP-MoO<sub>3</sub> at the fully discharged state. Scale bars: 50 nm.

The interesting H<sub>3</sub>O<sup>+</sup>-intercalation behavior of WP-MoO<sub>3</sub> is further supported by the Energy-dispersive X-ray spectroscopy (EDX) elemental mapping analysis. The fully discharged  $\alpha$ -MoO<sub>3</sub> presents the even distribution of Mo and Zn over the sample (Figure 3e & S12), while Zn distribution is barely observed in WP-MoO<sub>3</sub> (Figure 3f & S12). The Zn/Mo atomic ratio of the discharged WP-MoO<sub>3</sub> is calculated to be 0.02, which contrasts with the high Zn/Mo atomic ratio of the discharged  $\alpha$ -MoO<sub>3</sub> (0.48). Moreover, the discharged  $\alpha$ -MoO<sub>3</sub> and WP-MoO<sub>3</sub> were annealed in air at 500°C and subjected for XRD measurements (Figure S13). Peaks refer to ZnMoO<sub>3</sub> are only detected for the discharged  $\alpha$ -MoO<sub>3</sub>, rather than for the discharged WP-MoO<sub>3</sub>. All these results identify the successful modulation of intercalating charge carriers for  $\alpha$ -MoO<sub>3</sub>, which brings the obtained WP-MoO<sub>3</sub> with exceptional selective H<sub>3</sub>O<sup>+</sup>-intercalation chemistry in a neutral electrolyte.

To understand the origin of the selective H<sub>3</sub>O<sup>+</sup>-intercalation behavior, the first three CV cycles of WP-MoO<sub>3</sub> electrode were recorded as shown in Figure 4a. In the initial cycle, only a small cathodic peak is observed during the discharge process, whereas the charge process displays four anodic peaks corresponding to the extraction of the pre-intercalated H<sub>3</sub>O<sup>+</sup> and H<sub>2</sub>O. Afterwards, WP-MoO<sub>3</sub> electrode presents almost the identical 2<sup>nd</sup> and 3<sup>rd</sup> cycles (Figure S14), which include three anodic peaks and three cathodic peaks. The cathodic peaks and anodic peaks can be assigned to the H<sub>3</sub>O<sup>+</sup> intercalation and de-intercalation, respectively. Based on the CV curves, the charge transfer number of the first two pairs of redox peaks is calculated to be 1.495 (close to 1.5), implying the conversion between Mo<sup>6+</sup> and Mo<sup>5+</sup>/Mo<sup>4+</sup> (1:1). Meanwhile, the charge transfer of the third pair of redox

peaks is 0.498 (close to 0.5), which refers to the conversion between Mo<sup>5+</sup>/Mo<sup>4+</sup> (1:1) and Mo<sup>4+</sup>. In this regard, the structure of WP-MoO<sub>3</sub> after the first cycle (denoted WP-MoO<sub>3</sub>-c) is of significance to induce the selective H<sub>3</sub>O<sup>+</sup> intercalation of WP-MoO<sub>3</sub> electrode. The Mo 3d XPS spectrum of WP-MoO<sub>3</sub>-c uncovers the existence of Mo<sup>5+</sup> in WP-MoO<sub>3</sub>-c (Figure S15), indicating that H<sup>+</sup> ions were not fully extracted from the lattice. This result is further supported by the observation of O<sub>3</sub> peak in the O 1s XPS of WP-MoO<sub>3</sub>-c (Figure 4b). The partial extraction of H<sub>2</sub>O was verified by the larger O1 peak intensity of WP-MoO<sub>3</sub>-c than that of  $\alpha$ -MoO<sub>3</sub>, as well as the TGA analysis (Figure S16). Moreover, the interlayer distance of WP-MoO<sub>3</sub>-c identified by the XRD peak position only slightly decreased from 7.6 Å for WP-MoO<sub>3</sub> to 7.4 Å, which remains to be significantly larger than that of  $\alpha$ -MoO<sub>3</sub> (6.7 Å) (Figure 4c & S17).

Based on the quantitative analysis of XPS results, the structures of  $\alpha$ -MoO<sub>3</sub> and WP-MoO<sub>3</sub>-c (Figure S18) were simulated with the density functional theory (DFT) method. In WP-MoO<sub>3</sub>-c, the bonding interaction between the residual intercalant (i.e., H<sub>2</sub>O and H<sub>3</sub>O<sup>+</sup>) and the terminal O of [MoO<sub>6</sub>] bilayers is proved by the formant arise at  $-10 \text{ eV}$  in projected density of states (PDOS) of terminal O-p orbital with H-s orbital (Figure 4d). This bonding interaction greatly influence the electron density of terminal O atoms, which agrees well with the O 1s XPS results of our samples. Additionally, WP-MoO<sub>3</sub> has a band gap of 0.04 eV, which is remarkably smaller than that of  $\alpha$ -MoO<sub>3</sub> (2.18 eV). The Fermi level of WP-MoO<sub>3</sub> also shifts to the conduction band, favoring the excitation of charge carriers to the conduction band and thus the improved electronic conductivity (Figure S19).<sup>[21]</sup>

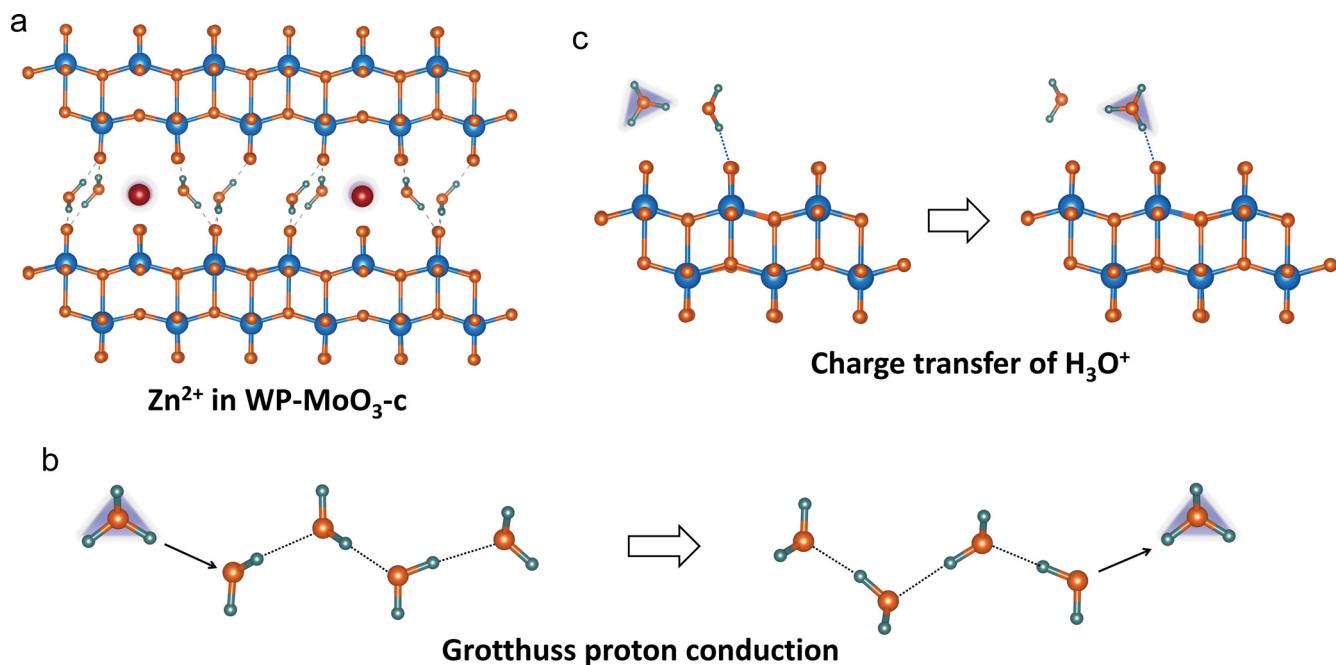


**Figure 4.** a) Initial three CV cycles of WP-MoO<sub>3</sub> at 0.7 mVs<sup>-1</sup>. b) O 1s XPS spectra of  $\alpha$ -MoO<sub>3</sub> and WP-MoO<sub>3</sub>-c. c) XRD patterns of  $\alpha$ -MoO<sub>3</sub>, WP-MoO<sub>3</sub>, and WP-MoO<sub>3</sub>-c. d) Simulated PDOS of  $\alpha$ -MoO<sub>3</sub> and WP-MoO<sub>3</sub>-c. Calculated e) Zn<sup>2+</sup>-intercalation energy and f) H<sub>3</sub>O<sup>+</sup>-intercalation energy for  $\alpha$ -MoO<sub>3</sub> and WP-MoO<sub>3</sub>-c.

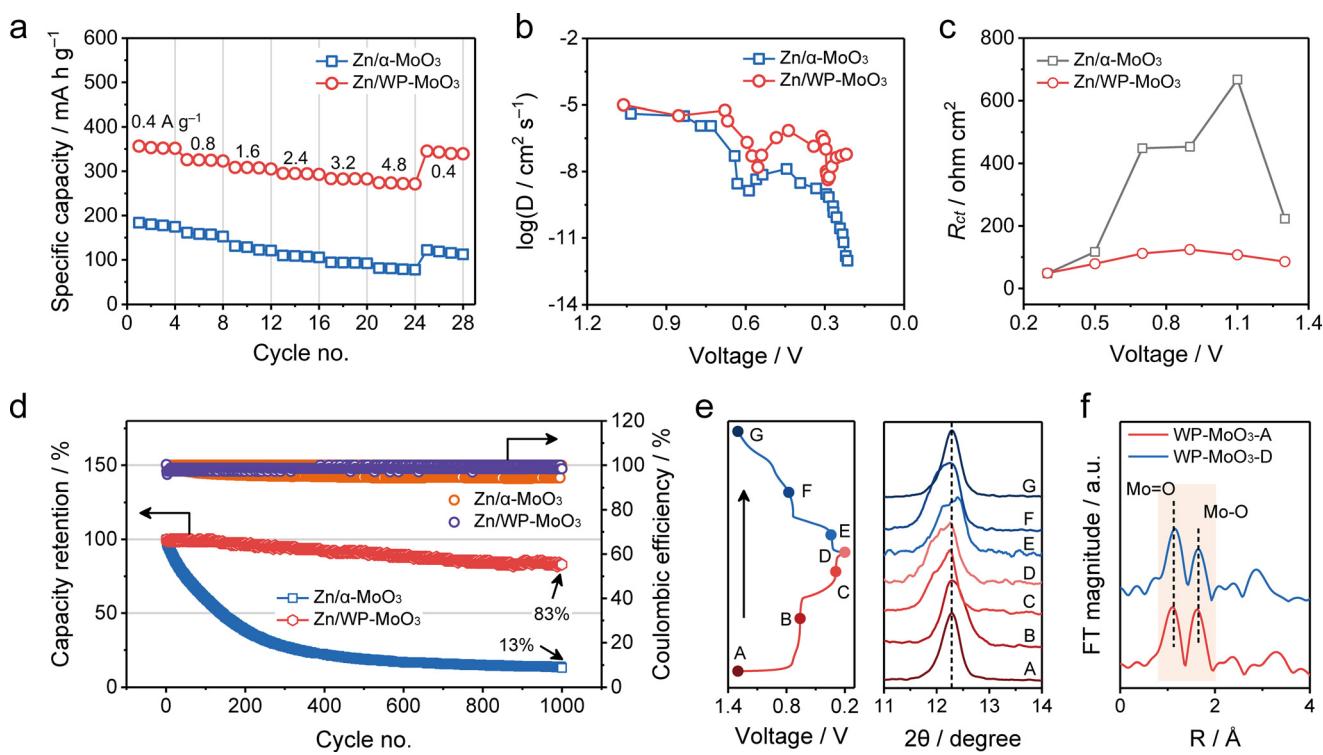
Subsequently, the Zn<sup>2+</sup>-intercalation energy (Figure 4e) and H<sub>3</sub>O<sup>+</sup>-intercalation energy (Figure 4f) were calculated for  $\alpha$ -MoO<sub>3</sub> and WP-MoO<sub>3</sub>-c. As expected, a substantially enlarged Zn-intercalation energy is uncovered for WP-MoO<sub>3</sub>-c ( $-0.11$  eV) in comparison to  $\alpha$ -MoO<sub>3</sub> ( $-1.49$  eV), while the H<sub>3</sub>O-intercalation energy of WP-MoO<sub>3</sub>-c ( $-3.93$  eV) is notably lower than that of  $\alpha$ -MoO<sub>3</sub> ( $-1.61$  eV). These results confirm the thermodynamically preferable H<sub>3</sub>O<sup>+</sup> intercalation into WP-MoO<sub>3</sub>-c. Besides, the residual H<sub>2</sub>O and H<sub>3</sub>O<sup>+</sup> located in the interlayer space and bonded with terminal O atoms of [MoO<sub>6</sub>] bilayers block the Zn<sup>2+</sup>-diffusion pathways in WP-MoO<sub>3</sub>-c (Figure 5a) and hinder the charge transfer through the interaction with [MoO<sub>6</sub>] bilayers (Figure S20 & 21). In the case of H<sub>3</sub>O<sup>+</sup> intercalation into WP-MoO<sub>3</sub>-c, charge carrier diffusion can be efficiently achieved through a well-established Grotthuss mechanism (Figure 5b), in which protons can be fast transported by “jumping” through water molecules.<sup>[22]</sup> Proton conductivities measurement (Figure S22) shows that WP-MoO<sub>3</sub>-c owns a proton conductivity value of  $4.3 \times 10^{-3}$  Scm<sup>-1</sup> at 318 K and 100% humidity, which is much higher than  $\alpha$ -MoO<sub>3</sub> ( $6.2 \times 10^{-5}$  Scm<sup>-1</sup>). Moreover, WP-MoO<sub>3</sub>-c also shows an activation energy ( $E_a$ ) of 0.28 eV, which suggests the Grotthuss conduction mechanism ( $E_a <$

0.4 eV).<sup>[23]</sup> H<sub>2</sub>O molecules between the interlayers serve as the proton transport intermedia, providing a hydrogen-bonding network for the high-kinetics charge carrier diffusion.<sup>[24]</sup> Moreover, the charge transfer from H<sub>3</sub>O<sup>+</sup> to WP-MoO<sub>3</sub>-c can be also through Grotthuss mechanism without breaking the hydrogen bonding interaction between H<sub>2</sub>O and the terminal O of [MoO<sub>6</sub>] bilayers (Figure 5c). Thereby, the unique van der Waals structure of WP-MoO<sub>3</sub>-c favors the selective H<sub>3</sub>O<sup>+</sup>-intercalation behavior both thermodynamically and kinetically.

The selective H<sub>3</sub>O<sup>+</sup> intercalation of WP-MoO<sub>3</sub> electrode motivated us to assess the electrochemical performance of ZMB devices assembled by coupling Zn anodes with WP-MoO<sub>3</sub> cathodes (denoted Zn/WP-MoO<sub>3</sub>). ZMB devices based on Zn<sup>2+</sup>-intercalation  $\alpha$ -MoO<sub>3</sub> cathodes were also constructed for comparison (denoted Zn/ $\alpha$ -MoO<sub>3</sub>). GCD curves at various current densities were collected to evaluate both devices (Figure S23). As shown in Figure 6a, Zn/ $\alpha$ -MoO<sub>3</sub> only exhibits a specific capacity of 184.0 mAh g<sup>-1</sup> at 0.4 A g<sup>-1</sup> (based on the cathode) and a capacity retention of 42.2% at 4.8 A g<sup>-1</sup>. By contrast, Zn/WP-MoO<sub>3</sub> delivers a much larger specific capacity (356.8 mAh g<sup>-1</sup> at 0.4 A g<sup>-1</sup>) and a greatly improved rate capability (77.5% from 0.4 to 4.8 A g<sup>-1</sup>). When



**Figure 5.** a) Schematic illustration of  $\text{Zn}^{2+}$  inserted  $\text{WP-MoO}_3\text{-c}$ . b) Schematic illustration of the Grotthuss proton-conduction mechanism. Protons are transported by rearranging bonds along a water chain. c) Schematic illustration of the charge transfer of  $\text{H}_3\text{O}^+$  via Grotthuss mechanism.



**Figure 6.** a) Rate performance, b) ion diffusion coefficient as a function of the voltage during the discharge, c)  $R_{\text{ct}}$  at different voltages, and d) cycling performance at  $3.2 \text{ A g}^{-1}$  of  $\text{Zn}/\alpha\text{-MoO}_3$  and  $\text{Zn}/\text{WP-MoO}_3$  devices. e) Ex situ XRD of  $\text{WP-MoO}_3$  during one discharge/charge cycle of  $\text{Zn}/\text{WP-MoO}_3$ , f) Radial distribution functions obtained from the  $k^2\chi(k)$  by Fourier transform of XANES of  $\text{WP-MoO}_3$  cathode at the charged ( $\text{WP-MoO}_3\text{-A}$ ) and discharged ( $\text{WP-MoO}_3\text{-D}$ ) states.

the current density returns to  $0.4 \text{ A g}^{-1}$  after the rate tests,  $\text{Zn}/\text{WP-MoO}_3$  can still show a specific capacity of  $345.6 \text{ mA h g}^{-1}$ . The influence of different amounts of intercalated  $\text{H}^+$  (from

0.25 to 2.0 per  $\text{MoO}_3$ ) on the electrochemical performance was also investigated by preparing a series of  $\text{WP-MoO}_3$  electrodes with different cut-off potentials (Figure S24). In

brief, the larger amount of intercalated  $\text{H}^+$  results in the higher specific capacity of the obtained Zn/WP-MoO<sub>3</sub> devices. Importantly, Zn/WP-MoO<sub>3</sub> delivers the maximum energy density of 198.0 Wh kg<sup>-1</sup> (based on the cathode) at a power density of 0.28 kW kg<sup>-1</sup>, as well as the peak power density of 6.7 kW kg<sup>-1</sup> at a high energy density of 104.5 Wh kg<sup>-1</sup>. These energy and power densities significantly outperform those of the recently reported ZMBs based on cathodes like MnO<sub>2</sub>,<sup>[3c]</sup> ZnMn<sub>2</sub>O<sub>4</sub>,<sup>[25]</sup> VS<sub>2</sub>,<sup>[26]</sup> ZnHCF,<sup>[27]</sup> and pyrene-4,5,9,10-tetraone<sup>[28]</sup> (Figure S25). It should be noticed that the energy contribution of plateau II ( $\approx 0.30$  V) is only about 16.8% in Zn/WP-MoO<sub>3</sub> (Figure S26), while most of the energy is contributed by plateau I ( $\approx 0.68$  V). The outstanding performance of Zn/WP-MoO<sub>3</sub> originates from H<sub>3</sub>O<sup>+</sup> charge carriers, which allow the high-kinetics diffusion and the efficient charge transfer. Figure 6b displays the calculated charge carrier diffusion coefficients of Zn/ $\alpha$ -MoO<sub>3</sub> and Zn/WP-MoO<sub>3</sub> devices based on a galvanostatic intermittent titration technique (GITT, Figure S27). Both devices exhibit low charge carrier diffusion coefficients at voltages associated with ion-intercalation stages. As expected, Zn/WP-MoO<sub>3</sub> with H<sub>3</sub>O<sup>+</sup> charge carriers achieves high diffusion coefficients of  $1.6 \times 10^{-8}$  cm<sup>2</sup> s<sup>-1</sup> at 0.57 V and  $4.4 \times 10^{-9}$  cm<sup>2</sup> s<sup>-1</sup> at 0.29 V, which are several orders of magnitude higher than those of Zn/ $\alpha$ -MoO<sub>3</sub> with Zn<sup>2+</sup> as the main charge carrier ( $1.2 \times 10^{-9}$  cm<sup>2</sup> s<sup>-1</sup> at 0.58 V and  $1.3 \times 10^{-12}$  cm<sup>2</sup> s<sup>-1</sup> at 0.22 V). In addition, the low charge transfer resistance ( $R_{ct}$ ) of Zn/WP-MoO<sub>3</sub> derived from electrochemical impedance spectroscopy (EIS, Figure S28) reflects the efficient charge transfer enabled by H<sub>3</sub>O<sup>+</sup> charge carriers. As shown in Figure 6c,  $R_{ct}$  of Zn/WP-MoO<sub>3</sub> ranges from 48.9 to 124.5  $\Omega$ , which contrasts with that of Zn/ $\alpha$ -MoO<sub>3</sub> (48.1 to 667.3  $\Omega$ ).

The cycling stability of Zn/ $\alpha$ -MoO<sub>3</sub> and Zn/WP-MoO<sub>3</sub> was evaluated at a current density of 3.2 A g<sup>-1</sup>. Zn/WP-MoO<sub>3</sub> presents impressive coulombic efficiencies of nearly 100%, which indicates its high charge/discharge reversibility. In contrast with the fast capacity decay of Zn/ $\alpha$ -MoO<sub>3</sub> (13% capacity retention after 1000 cycles), Zn/WP-MoO<sub>3</sub> can maintain 83% of the initial capacity after 1000 cycles (Figure 6d). The outstanding cycling performance of Zn/WP-MoO<sub>3</sub> is assigned to the negligible structure distortion of WP-MoO<sub>3</sub> during the repeated H<sub>3</sub>O<sup>+</sup> intercalation/extraction. In the ex-situ XRD tests of WP-MoO<sub>3</sub> cathode during one discharge/charge cycle of Zn/WP-MoO<sub>3</sub>, the peak position located at 12.3° of WP-MoO<sub>3</sub> only experiences slight shift, indicating the little volume expansion/shrinkage of WP-MoO<sub>3</sub> in direction which is perpendicular to interlayer (Figure 6e). Moreover, Mo-K edge XANES (Figure S29) and corresponding Fourier transforms of the Mo K-edge  $k^2\chi(k)$  spectra (Figure 6f) confirm that both Mo=O and Mo—O bonds of WP-MoO<sub>3</sub> cathode have no change in bonding length at the fully charged and discharged stages. All these results imply that the selective H<sub>3</sub>O<sup>+</sup>-intercalation chemistry is able to empower electrodes with large specific capacity, high charge-storage kinetics and long cycle life.

## Conclusion

In summary, we have uncovered an exceptional selective H<sub>3</sub>O<sup>+</sup>-intercalation chemistry for WP-MoO<sub>3</sub> electrode in a neutral ZnCl<sub>2</sub> electrolyte. The interesting charge-carrier-selection behavior of WP-MoO<sub>3</sub> originated from the interlayer species (i.e., H<sub>2</sub>O, H<sub>3</sub>O<sup>+</sup>) of WP-MoO<sub>3</sub>, which allowed the fast-kinetics transport and charge transfer of H<sub>3</sub>O<sup>+</sup> while blocking Zn<sup>2+</sup> intercalation. This study provides a novel charge-carrier-modulation concept through the strategic van der Waals structure engineering, which opens a promising prospect for developing high-kinetics and long-life battery technologies.

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## Conflict of interest

The authors declare no conflict of interest.

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